

# SUMMARY OF IONESCU-SCHLAG AND SOME OTHER THINGS

## 1. SUMMARY OF IONESCU-SCHLAG [7]

Recall some notation.

- $B = \{f : \mathbb{R}^d \rightarrow \mathbb{C} : \|f\|_B^2 := \sum_{j \geq 0} 2^j \int_{|x| \in [2^{j-1}, 2^j]} |f|^2 < \infty\}$  and  $B^* = \{g : \mathbb{R}^d \rightarrow \mathbb{C} : \|g\|_{B^*}^2 := \sup_{j \geq 0} 2^{-j} \int_{|x| \in [2^{j-1}, 2^j]} |g|^2 < \infty\}$  the classical Agmon-Hörmander spaces. Note  $B \hookrightarrow L^2 \hookrightarrow B^*$ .
- $S_\alpha : \mathcal{S}'(\mathbb{R}^d) \rightarrow \mathcal{S}'(\mathbb{R}^d)$  with  $S_\alpha f = (1 - \Delta)^{\alpha/2} f$  and  $\alpha \in \mathbb{C}$ . (For  $\alpha \in \mathbb{R}$  this corresponds to fractional differentiation/integration).
- $W^{\alpha,p} = \{u \in \mathcal{S}'(\mathbb{R}^d) : S_\alpha u \in L^p\}$  with  $\|u\|_{W^{\alpha,p}} = \|S_\alpha u\|_{L^p}$  denote the usual Sobolev spaces. If  $p = 2$ , then  $W^{\alpha,2} \equiv H^\alpha$ .
- For  $p_d = (2d + 2)/(d + 3)$  the main Banach spaces<sup>[1]</sup> here are

↑ reg. + schnell abf.

$$X = \underbrace{W^{-1/(d+1), p_d}} + S_1(B), \quad \|f\|_X = \inf_{f=f_1+f_2} \|S_{-1/(d+1)} f_1\|_{L^{p_d}} + \|S_{-1} f_2\|_B$$

reg. + langsam abf.

$$X^* = W^{1/(d+1), p'_d} \cap S_{-1}(B^*), \quad \|g\|_{X^*} = \max\{\|S_{1/(d+1)} g\|_{L^{p'_d}}, \|S_1 g\|_{B^*}\}$$

Note that  $X^* \subseteq H^1_{loc}$  and  $X \hookrightarrow H^{-1}$  and  $H^1 \hookrightarrow X^*$ .

**Theorem 1.1** (Combined LAP). Let  $\delta \in (0, 1]$ , then

$$R_0(z) = (-\Delta - z)^{-1}$$

$L^p$   
 $L^p \subset L^q$   
 $p < q$

$$\sup_{|\lambda| \in (\delta, \delta^{-1})} \sup_{\varepsilon \in (-1, 1) \setminus 0} \|R_0(\lambda \pm i\varepsilon)\|_{X \rightarrow X^*} \lesssim \delta^{-1}$$



The following is a weighted version thereof and will be used to establish<sup>[2]</sup>

- discreteness of point spectrum in  $\mathbb{R} \setminus 0$  and
- rapid decay of eigenfunctions.

For  $N \geq 0$  and  $\gamma \in (0, 1]$ , let

$$\mu_{N,\gamma}(x) = \frac{\langle x \rangle^{2N}}{\langle \sqrt{\gamma} x \rangle^{2N}} \quad \xrightarrow{\gamma \rightarrow 0} \langle x \rangle^{2N}$$

which equals  $\langle x \rangle^{2N}$  for  $\gamma \rightarrow 0$  and 1 for  $\gamma = 1$ .

**Theorem 1.2** (Weighted combined LAP). Let  $\delta \in (0, 1]$ , then

$$\|\mu_{N,\gamma} u\|_{X^*} \lesssim_{N,\delta} \|\mu_{N,\gamma} (-\Delta - \lambda) u\|_X, \quad |\lambda| \in (\delta, \delta^{-1})$$

whenever  $u \in X^*$  satisfies the mild decay condition (cf. Lemma 4.3)

$$\lim_{R \rightarrow \infty} R^{-1} \int_{R \leq |x| \leq 2R} |u|^2 = 0.$$

<sup>1</sup> $X$ : fast decay, low regularity —  $X^*$ : slow decay, good regularity.

<sup>2</sup>Compare with [3] Theorem 14.2.4].

$$H' \supset X^* \subset H'_{loc}$$

**Definition (Admissible potentials).** Let  $\mathfrak{B}(X^*, X)$  denote the space of linear bounded operators from  $X^*$  to  $X$ . Then  $V$  is said to be admissible if

- (1)  $V \in \mathfrak{B}(X^*, X)$  and  $(\psi, V\varphi) = \overline{(\varphi, V\psi)}$  for any  $\psi, \varphi \in \mathcal{S}(\mathbb{R}^d)$ .
- (2) For all  $\varepsilon > 0$  and  $N \geq 0$  there exist  $A_{N,\varepsilon}, R_{N,\varepsilon} \geq 1$  such that

$$V: X^* \rightarrow X$$

$$\|\mu_{N,\gamma} V u\|_X \leq \varepsilon \|\mu_{N,\gamma} u\|_{X^*} + A_{N,\varepsilon} \|u \mathbf{1}_{|x| \leq R_{N,\varepsilon}}\|_{L^2}, \quad u \in X^*, \gamma \in (0, 1].$$

(Observe what happens if  $u$  solves  $-\Delta + V u = \lambda u$  in view of Theorem 1.2).

- (3) There exist  $J \in \mathbb{N}_{>0}$  and operators  $A_j, B_j \in \mathfrak{B}(X^*, L^2)$  ( $j = 1, 2, \dots, J$ ) such

Kato-  
Orelshel'ny

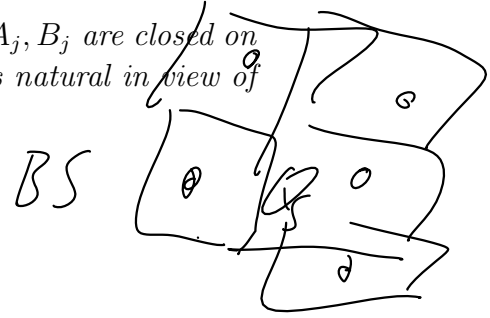
that

$$V = B^* A \quad (\varphi, V\psi) = \sum_{j=1}^J (B_j \varphi, A_j \psi), \quad \varphi, \psi \in X^*.$$

Moreover, considered as (unbounded) operators on  $L^2$ , the  $A_j, B_j$  are closed on some domains satisfying  $\mathcal{D}(A_j), \mathcal{D}(B_j) \supseteq H^1(\mathbb{R}^d)$  (which is natural in view of  $H^1 \hookrightarrow X^*$ .) (Recall Kato smoothing theory in Section 2.)

For instance  $V \in \ell^{(d+1)/2} L^{d/2}$  are admissible where

$$\frac{d}{2}, \quad \frac{d+1}{2} \text{ vs } \frac{d}{2} \quad \|V\|_{\ell^{(d+1)/2} L^{d/2}} = \left[ \sum_s \|V\|_{L^{\frac{d}{2}}(Q_s)}^{\frac{d+1}{2}} \right]^{\frac{2}{d+1}}$$



and  $\{Q_s\}_{s \in \mathbb{Z}^d}$  is a collection of axis-parallel unit cubes such that  $\mathbb{R}^d = \bigcup_s Q_s$ . The exponent  $d/2$  indicates the local integrability which is optimal to realize  $-\Delta + V$  self-adjointly, whereas the exponent  $(d+1)/2$  indicates its decay. (Note that  $2/(d+1) = 1/p_d - 1/p'_d$ .)

**Theorem 1.3** (Agmon–Hörmander–Kato–Kuroda for admissible  $V$ ). Let  $V$  be admissible, then

- (1)  $H \equiv -\Delta + V$  defines a self-adjoint operator on  $\mathcal{D}(H) = \{u \in H^1 : Hu \in L^2\}$ . Moreover,  $H \geq -c$  for some  $c \in \mathbb{R}$ .
- (2)  $\mathcal{E} := \sigma_{pp} \setminus 0$  is discrete in  $\mathbb{R} \setminus 0$ , i.e.,  $\mathcal{E} \cap I$  is finite for any compact  $I \subseteq \mathbb{R} \setminus 0$ . Moreover, each eigenvalue has at most finite multiplicity.
- (3) The eigenfunctions  $u$  of  $H$  decay rapidly, i.e., for each  $N \geq 0$ , one has that  $\langle x \rangle^N u \in H^1(\mathbb{R}^d)$ .
- (4) We have a LAP for  $H$ , i.e., for compact  $I \subseteq (\mathbb{R} \setminus 0) \setminus \mathcal{E}$ , we have

$$\sup_{\lambda \in I, \varepsilon \in (-1, 1) \setminus 0} \|(-\Delta + V - (\lambda \pm i\varepsilon))^{-1}\|_{X \rightarrow X^*} \lesssim_{V, I} 1.$$

In particular,  $\sigma(H) \cap I = \sigma_{ac}(H) \cap I$ , i.e., the spectrum of  $H$  on  $I$  is purely ac.

- (5)  $\sigma_{sc}(H) = 0$  and  $\sigma_{ac}(H) = \sigma_{ac}(-\Delta) = [0, \infty)$ .

- (6) The generalized wave operators  $\Omega_{\pm}(H, H_0) := s - \lim_{t \mp \infty} e^{itH} e^{-itH_0} P_{ac}^{(0)}$  and  $\Omega_{\pm}(H_0, H) := s - \lim_{t \mp \infty} e^{itH_0} e^{-itH} P_{ac}^{(V)}$  exist and are complete.

**Remark** (Embedded eigenvalues). In principle the theorem does not rule out point spectrum in  $[0, \infty)$ . However, a deep result by Koch and Tataru [8 Theorem 3] actually

$$\underline{(H-z)^{-1}} = \underbrace{\left( \mathbb{1} + R_0(z)V \right)^{-1}}_{\text{Lippmann-Schwinger-}} \underline{R_0(z)}$$

4.1

a)  $\Gamma: (0, \infty) \rightarrow \mathcal{L}(X, X^*)$  ist analytisch  
 $z \mapsto R_0(z)$

$\rightarrow$  b)  $\forall \lambda \in (0, \infty)$ , existieren Operatoren  $R_0(\lambda + i0)$

$$\|R_0(\lambda + i0)\|_{X \rightarrow X^*} \leq 1$$

$$(R_0(\lambda + i\varepsilon)g, \phi) \xrightarrow{\varepsilon \rightarrow 0} (R_0(\lambda + i0)g, \phi), \quad g \in X, \phi \in \mathcal{S}$$

$$\| \mathbb{1}_{X \times X} [R_0(\lambda + i\varepsilon) - R_0(\lambda + i0)]g \| \xrightarrow{\varepsilon \rightarrow 0} 0 \quad g \in X.$$

$$c) (-\Delta - \lambda - i\varepsilon)R_0(\lambda + i\varepsilon)g \stackrel{X}{=} g \quad \text{im } \mathcal{S}'\text{-Sinn} \\ g \in X.$$

Pf

a)  $\checkmark$

c)  $\lambda < 0$ , oder  $\varepsilon \neq 0$ .  $\checkmark$   
 $\lambda > 0$ ,  $\varepsilon = 0$  folgt aus b)

$$b) f \in \mathcal{S} \quad (-\Delta - z)^{-1}f(x) = T_z * f(x)$$

$$|(-\Delta - z)^{-1}(x)| \lesssim \underbrace{|x|^{-(d-2)}}_{\text{für } \operatorname{Im} z} \mathbb{1}_{|x| < 1} + \underbrace{|x|^{-\frac{d-1}{2}}}_{\frac{1}{\sqrt{\delta}}(x)} \mathbb{1}_{|x| > 1}$$

$$T_{d+i\epsilon} * f(x) \xrightarrow{\epsilon \rightarrow 0} T_{d+i0} * f(x) \quad (\text{kompakt})$$

$$\|R_0(d+i0) f\|_{X^*} \leq \lim_{\epsilon \rightarrow 0} \|T_{d+i\epsilon} * f\|_{X^*} \\ \stackrel{!}{\leq} \|f\|_X$$

Prinzip d.  $\Gamma$   $\rightarrow$  FAP  $\Gamma$  per Stetigkeit auf Fkte in  $X$   
 glm. Beschr.  $\rightarrow$  hochziehen

$$(R_0(d+i\epsilon)g, \phi) \rightarrow (R_0(d+i0)g, \phi) \quad g \in X \\ g \in S \text{ mit maj. Kon.} \quad \hookrightarrow$$

4.2  $V: X^k \rightarrow X$  kompakt

Pf  $V = V_1 + V_2$   
 $\underbrace{V_1}_{C_c} \quad \underbrace{V_2}_{X^k \rightarrow X} < \epsilon$   
 $\hookrightarrow$  Rellich

4.3  $g \in X$

$$\text{Im}(g, R_0(d+i0)g) = c_1 \int_{|\xi|=\sqrt{\lambda}} |\hat{g}(\xi)|^2 d\sigma(\xi)$$

$\phi \in C_c^\infty$   
 (Fall  $X = \text{say} - \text{Transf}$ )  
 $\hookrightarrow \langle X, S \rangle$

$$\int_{\mathbb{R}} |(R_0(d+i0)g)(x)|^2 \phi\left(\frac{x}{R}\right) \frac{dx}{R} \xrightarrow{R \rightarrow \infty} c_2(\phi) \int_{|\xi|=\sqrt{\lambda}} |\hat{g}(\xi)|^2 d\sigma$$

$d\sigma(\xi) = \frac{d\mathcal{L}(\xi)}{2|\xi|}$

$\phi \int_{-\infty}^{\infty} \phi(x+t\omega) dt \hookrightarrow \forall x \in \mathbb{R}^n, \omega \in S^{n-1}$

Pf  $\lambda \in \mathbb{S} \checkmark$  ( $\rightarrow$  Hörmander I, II)

$\lambda \in \mathbb{X} \rightarrow$  mit TS oder Agmon-Hörmander

$$\text{L.4} \quad \underline{\underline{(H-z)^{-1}}} = \underline{\underline{(I + R_0(z)V)^{-1}}} \underline{\underline{R_0(z)}}$$

$$\xi = \sigma_{pp}(H) \setminus 0$$

$$\tilde{\xi} = \{ \lambda \in \mathbb{R} : \exists f \in X^* : (I_{X^*} + R_0(\lambda+i0)V)f = 0 \}$$

$$\rightarrow \mathcal{F}_\lambda = \{ f \in X^* : (I_{X^*} + R_0(\lambda+i0)V)f = 0 \}$$

Für  $\tilde{\xi} = \xi$ ,  $\tilde{\xi}$  dicht.

$$\text{L.4} \quad \text{Sei } u \in \mathcal{F}_\lambda \Rightarrow \| \langle x \rangle^\nu u \|_{X^*} \leq \| u \|_{X^*}.$$

$$\text{Pf} \quad \| \langle x \rangle^\nu u \|_{X^*} \stackrel{1.2}{\leq} \| \langle x \rangle^\nu (I + R_0(\lambda+i0)V)u \|_{X^*} \\ = \| \langle x \rangle^\nu u - R_0(\lambda+i0)Vu \|_{X^*}$$

$$\stackrel{1.1}{=} \| \langle x \rangle^\nu Vu \|_{X^*}$$

$$\forall \epsilon > 0 \quad \exists \delta \in \| \langle x \rangle^\nu u \|_{X^*} + \delta \leq \| u \|_{X^*}$$

$$\lim_{R \rightarrow \infty} \int_{\mathbb{R}^n} |u|^2 \frac{dx}{R} \stackrel{1.3}{=} \int_{|\xi|=\frac{1}{R}} d\sigma(\xi) | \widehat{Vu}(\xi) |^2$$

$$= \int_{\mathbb{R}^n} \underbrace{-R_0 V u}_{= -u}$$

$$\stackrel{1.3}{=} \int_{\mathbb{R}^n} (Vu, R_0(\lambda+i0)Vu) = -u$$

$$= 0 \quad \square$$

4.5  $\tilde{\varepsilon} \in \varepsilon$ ,  $\tilde{\varepsilon}$  diskret  $(H' \hookrightarrow) X^* \subseteq H'_{loc}$

Pf  $(\mathbb{1}_{X^*} + \mathcal{R}_0(\lambda + i0)V)^{-1} u = 0$ ,  $u \in F_\lambda$

4.1  $(-\Delta - \lambda) + Vu = 0$ ;  $u \in H'$  ✓

↳ h.h. u fällt schnell.

Ang  $F \propto$  EW  $\lambda_n$  mit  $u_n$  EV.

Beh:  $\|u_n\|_{H'} = 1$

Beh  $F$  konv. T.f. von  $u_n$  in  $H'$

$$u_n = \mathcal{R}_0(-1) \left[ (\lambda_n + 1 - V) u_n \right]$$

$$= (\lambda_n + 1) \underbrace{\mathcal{R}_0(-1) \langle X \rangle^{-2} \langle X \rangle^2}_{\text{komp. in } H'} u_n - \underbrace{\mathcal{R}_0(-1)V}_{\text{komp. in } H'} u_n$$

$u_n \xrightarrow{H'} u_\infty$   $\|u_\infty\|_{H'} = 1$  ✓

$(u_n, u_m) = \delta_{nm}$ .  $\rightarrow u_\infty = 0$

↳

4.6  $\|(\mathbb{1}_{X^*} + \underbrace{\mathcal{R}_0(\lambda + i\varepsilon)V}_{\text{komp. ht}})^{-1}\|_{X^*, X^*} \approx 1$ ,  $\varepsilon \in [0, 1]$   
 $\lambda \in \mathbb{R} \setminus 0$

→ analytisch Fredholm bis auf einer diskrete iZ in  $\mathbb{C}$  invertierbar.

$\varepsilon = 0$   $\tilde{\varepsilon}$   $\varepsilon > 0$

Pf 1.3. c)  $Hu = \lambda u, \quad u \in \mathcal{D}(H)$

$$\| \langle x \rangle^N u \|_{X^*} \stackrel{1.2}{\leq} \| \langle x \rangle^N \underbrace{(\Delta + \lambda) u}_{=Vu} \|_X$$

$$\leq \epsilon \| \langle x \rangle^N u \|_{X^*} + a_\epsilon \| u \|_{X^*} \quad \checkmark$$

1)  $\sigma_{pp}(0)$  diskret

$$\epsilon \subseteq \tilde{\epsilon}$$

~~$\tilde{\epsilon} \subseteq \epsilon$~~   $\tilde{\epsilon} \subseteq \epsilon \quad \checkmark$   
 $\uparrow$  diskret

$$(-\Delta - b)u + Vu = 0, \quad u \in H'$$

$$\underbrace{R_0(b)(-\Delta - \lambda)u}_{\stackrel{?}{=} u} + R_0(b)Vu = 0$$

$$\underbrace{R_0(b)(-\Delta - \lambda)u}_{\in H' \rightarrow X} = u' \in X^* \quad / \quad -\Delta - \lambda$$

$$\text{Lnl} \quad (-\Delta - \lambda)(u - u') = 0;$$

$$\| u - u' \|_{X^*} \stackrel{(1.2)}{\leq} \| (\Delta + b)(u - u') \|_X = 0$$

$$u \quad \checkmark \quad u \in \mathcal{D}(H)$$

$$\lim_{k \rightarrow \infty} \int_{|x| < R} |u'(x)|^2 \frac{dx}{R} \stackrel{4.3}{=} - \int_{|x|=R} \underbrace{((\Delta + b)u(x))}_{=0 \text{ pl\u00fctweise}} |^2 d\sigma(x)$$

$\Rightarrow \hat{u} \in C_0^\infty$  f\u00e4hrt schnell

$$(H-z)^{-1} = \underbrace{(1 + R_0(z)V)^{-1}}_{\text{kompakt}} \underbrace{R_0(z)}_{\text{LTPP } \checkmark}$$

$$\rightarrow \| (H-z)^{-1} \|_{X, X^*} \lesssim 1, \quad z = d + i\epsilon$$

$\epsilon > 0$   
~~...~~

$$\Rightarrow \sigma_{sc} = \emptyset$$

$$d \in \mathbb{R} \setminus \{0\}$$

$$\sigma_{ac} = [0, \infty)$$

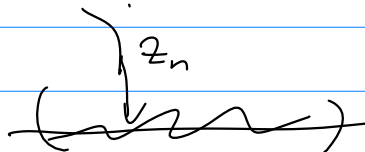
$$(\varphi, dE(\lambda)\varphi)$$

Mapt  $\int \frac{d\rho(\lambda)}{d-z} = F(z)$

$(\varphi, \int_{\Gamma} R(d+i\epsilon) \varphi)$   $\hookrightarrow$  Borel-Transf.  $\infty$   $\hookrightarrow$  von Neumann



$$\| (T-z)^{-1} \| > \frac{1}{\text{dist}(z, \sigma(T))}$$



$\downarrow (X \rightarrow X^*), -d, -d+i\epsilon \dots$   $\downarrow$   $\text{glm. in } \text{Im } z$

$$\| (T - d + i\epsilon)^{-1} \|_{X \rightarrow X^*} \lesssim 1$$



says that for  $V \in \ell^{(d+1)/2} L^{d/2}$  there are no embedded eigenvalues. This is sharp in view of the sequence of counterexamples<sup>3</sup> of Ionescu–Jerison [6] satisfying

$$\text{Ionescu-Jerison } L^p(\mathbb{R}^d) \ni |V_n(x)| \lesssim \frac{1}{n + \underbrace{(|x'|^2)^{3/2}}_{\text{underline}} + \underbrace{|x_d|}_{\text{underline}}}, \quad p > \frac{d+1}{2} \quad (1.1)$$

and  $\lim_{n \rightarrow \infty} \|V_n\|_{L^p(\mathbb{R}^d)} = 0$ <sup>4</sup>. Cuenin [1] recently related that counterexample to Knapp's counterexample in  $\|\hat{f}|_S\|_{L^q} \lesssim \|f\|_{L^p(\mathbb{R}^d)}$  which gave  $1/q \geq (d+1)/(p'(d-1))$ . Let  $a \in \mathbb{R}^d$ ,  $\delta \ll 1$ , and take  $f(x) = \chi_{\{|x'-a'| < \delta^{-1}, |x_d-a_d| < \delta^{-2}\}}(x)$ . Then  $\hat{f}$  is morally supported on the dual rectangle  $R_\delta^* = \{|\xi'| < \delta, |\xi_d| < \delta^2\}$ <sup>5</sup>. Now pave  $\mathbb{R}^d \setminus B_0(1)$  with finitely overlapping rectangles of dimension  $(2^j)^{d-1} \times 4^j$ , let  $\mathbf{1}_{\mathbb{R}^d \setminus B_0(1)}(x) = \sum_{j \geq 1} \chi_j(x)$ , and  $u(x) = \sum_{j \geq 1} 4^{-Nj} \chi_j(x)$  (being a superposition of Knapp examples). Then  $|u(x)| \sim (|x'|^2 + |x_d|)^{-N}$  and, since  $\Psi$ DOs do not move the support too much,

$$T(D)u(x) = \sum_{j \geq 1} 4^{-Nj} (T|_{R_j^*} + O(4^{-j})) \chi_j(x) \sim \lambda u(x) + O((|x'|^2 + |x_d|)^{-N-1})$$

where we expanded  $T(\xi) = \lambda + (\partial_d T(0))\xi_d + O(|\xi|^2)$  around the aspired eigenvalue  $\lambda \in \mathbb{C}$ . Thus,  $V = -(T(D) - \lambda)u/u$  is smooth and satisfies pointwise bound in (1.1).

**Lemma 4.1** (Analyticity of free resolvent and existence of boundary values).

(1) The map

$$\begin{aligned} \mathbb{C} \setminus [0, \infty) &\rightarrow \mathfrak{B}(X, X^*) \\ z &\mapsto R_0(z) \end{aligned}$$

is analytic.

(2) For any  $\lambda \in (0, \infty)$  there are operators  $R_0(\lambda \pm i0) \in \mathfrak{B}(X, X^*)$  such that

$$\|R_0(\lambda \pm i0)\|_{X \rightarrow X^*} \lesssim_\delta 1, \quad \lambda \in (\delta, \delta^{-1}), \delta \in (0, 1).$$

Moreover, for any sequences  $(\lambda_n)_{n \in \mathbb{N}} \subseteq (0, \infty)$  and  $(\varepsilon_n)_{n \in \mathbb{N}} \subseteq [0, \infty)$  with  $\lambda_n \rightarrow \lambda$  and  $\varepsilon_n \rightarrow 0$ , we have

$$\begin{aligned} (R_0(\lambda_n \pm i\varepsilon_n)f, \varphi) &\rightarrow (R_0(\lambda \pm i0)f, \varphi), \quad f \in X, \varphi \in \mathcal{S} \\ \|\mathbf{1}_{|x| \leq R} [R_0(\lambda_n \pm i\varepsilon_n) - R_0(\lambda \pm i0)]f\| &\rightarrow 0, \quad f \in X, R \geq 1. \end{aligned}$$

(3) For  $\lambda \in \mathbb{R} \setminus 0$  and  $\varepsilon \geq 0$ , we have

$$(-\Delta - (\lambda \pm i\varepsilon))R_0(\lambda \pm i\varepsilon)g = g, \quad g \in X$$

in distributional sense, i.e., whenever the inner product with  $\mathcal{S}$  functions is taken.

**Lemma 4.2.** If  $V$  is admissible, then  $V : X^* \rightarrow X$  is compact.

<sup>3</sup>Their construction is analogous to that of Wigner and von Neumann [10].

<sup>4</sup>This is interesting in view of the Laptev–Safronov conjecture, cf. Frank–Simon [2].

<sup>5</sup>A Knapp example living at the south pole  $\xi_d \sim 0$ .

The following establishes the link between Fourier restriction and the existence of boundary values of the free resolvent. See also Hörmander [3] Theorem 14.2.2 and Corollary 14.3.10] and Section 3].

**Lemma 4.3.** *Let  $\Phi \in C_c^\infty(\mathbb{R}^d)$  with  $\Phi(0) = 1$  and  $\Phi(x) = 0$  whenever  $|x| \geq 1$ . Then for any  $\lambda > 0$  and  $g \in X$ , we have*

$$\begin{aligned} \Im(g, R_0(\lambda \pm i0)g) &= c_1(\lambda, \pm) \int_{\sqrt{\lambda}\mathbb{S}^{d-1}} |\hat{g}(\xi)|^2 d\sigma(\xi) \\ \lim_{R \rightarrow \infty} \int_{\mathbb{R}^d} |(R_0(\lambda \pm i0)g)(x)|^2 \Phi\left(\frac{x}{R}\right) \frac{dx}{R} &= c_2(\lambda, \Phi, \pm) \int_{\sqrt{\lambda}\mathbb{S}^{d-1}} |\hat{g}(\xi)|^2 d\sigma(\xi) \end{aligned}$$

where  $d\sigma(\xi)$  denotes the ‘‘canonical’’ surface measure, i.e.,  $d\sigma(\xi) = d\Sigma(\xi)/(2|\xi|)$  where  $d\Sigma(\xi)$  denotes the euclidean (Lebesgue) surface measure.

Recall the Leray measure  $P^*\delta_0 = d\Sigma(\xi)/|\nabla P(\xi)|$  whenever  $P : \mathbb{R}^d \rightarrow \mathbb{R}$  with  $\nabla P \neq 0$ ,  $\delta_0 \in \mathcal{E}'(\mathbb{R}^d)$  is the usual  $d$ -dimensional Dirac distribution at the origin, and  $d\Sigma$  is the euclidean surface measure on  $\{\xi \in \mathbb{R}^d : P(\xi) = 0\}$ , see also Hörmander [4] Theorem 6.1.5].

Next, the relation between eigenfunctions of  $H$  and solutions to  $(\mathbf{1}_{X^*} + R_0(\lambda \pm i0)V)f = 0$  (the ‘‘Lippmann–Schwinger equation’’) is discussed. Let

$$\begin{aligned} \tilde{\mathcal{E}}^\pm &= \{\lambda \in \mathbb{R} \setminus 0 : \exists f \in X^* \setminus 0 \text{ s.t. } (\mathbf{1}_{X^*} + R_0(\lambda \pm i0)V)f = 0\} = \tilde{\mathcal{E}} \\ \mathcal{F}_\lambda^\pm &= \{f \in X^* : (\mathbf{1}_{X^*} + R_0(\lambda \pm i0)V)f = 0\}. \end{aligned}$$

Our goal is to show  $\tilde{\mathcal{E}} = \mathcal{E} := \sigma_{pp}(H) \setminus 0$ .

**Lemma 4.4** (Rapid decay of solutions to ‘‘Lippmann–Schwinger’’). *Let  $\lambda \in \tilde{\mathcal{E}}$  with corresponding  $f \in F_\lambda^\pm$ . Then for any  $N \geq 0$ , we have*

$$\| \langle x \rangle^{2N} f \|_{X^*} \lesssim_{N,V,\lambda} \|f\|_{X^*}.$$

Let  $\mathfrak{H}_\lambda = \{u \in \mathcal{D}(H) : Hu = \lambda u\}$  denote the vector space of eigenfunctions of  $H$ .

**Lemma 4.5** (Preliminary relationship between  $\mathcal{E}$  and  $\tilde{\mathcal{E}}$ ).

- (1) For any  $\lambda \in \mathbb{R} \setminus 0$  we have  $\mathcal{F}_\lambda^+ \cup \mathcal{F}_\lambda^- \subseteq \mathfrak{H}_\lambda$  and so  $\tilde{\mathcal{E}} \subseteq \mathcal{E}$ .
- (2)  $\tilde{\mathcal{E}}$  is discrete in  $\mathbb{R} \setminus 0$ , i.e.,  $I \cap \tilde{\mathcal{E}}$  is finite for any compact  $I \subseteq \mathbb{R} \setminus 0$ .
- (3) For any  $\lambda \in \mathbb{R} \setminus 0$  the vector spaces  $\mathcal{F}_\lambda^\pm$  are finite-dimensional.

**Lemma 4.6** (Uniform invertibility of  $\mathbf{1}_{X^*} + R_0(\lambda \pm i\varepsilon)V$  away from  $\tilde{\mathcal{E}}$ ). *For any  $\lambda \in (\mathbb{R} \setminus 0) \setminus \tilde{\mathcal{E}}$  the operators  $\mathbf{1}_{X^*} + R_0(\lambda \pm i\varepsilon)V$  are invertible on  $X^*$  with*

$$\sup_{\lambda \in I} \sup_{\varepsilon \in [0,1]} \|(\mathbf{1}_{X^*} + R_0(\lambda \pm i\varepsilon)V)^{-1}\|_{X^* \rightarrow X} \lesssim_I 1, \quad I \subseteq (\mathbb{R} \setminus 0) \setminus \tilde{\mathcal{E}} \text{ compact}.$$

**Lemma 5.1** (Resolvent formula for  $(H - z)^{-1}$  away from  $\mathbb{R}$ ). *For any  $\lambda \in \mathbb{R}$  and  $\varepsilon \in \mathbb{R} \setminus 0$ , the operator*

$$\tilde{R}_V(\lambda + i\varepsilon) := (\mathbf{1}_{H^1} + R_0(\lambda \pm i\varepsilon)V)^{-1}R_0(\lambda + i\varepsilon) : L^2 \rightarrow \mathcal{D}(H).$$

is well defined and bounded. Moreover, it is a right inverse, i.e.,

$$[H - (\lambda + i\varepsilon)]\tilde{R}_V(\lambda + i\varepsilon) = \mathbf{1}_{L^2}.$$

2. ELEMENTS OF KATO SMOOTHING THEORY

We recall some classic facts, cf. [9 Section XIII.7].

**Definition** (*H*-smoothness). Let *A* be closed and *H* self-adjoint. Then *A* is *H*-smooth if and only if

$$\begin{aligned} & \sup_{\varepsilon > 0, \|\varphi\|=1} \int_{\mathbb{R}} (\|AR(\lambda + i\varepsilon)\varphi\|^2 + \|AR(\lambda - i\varepsilon)\varphi\|^2) d\lambda < \infty \\ \text{or} & \sup_{\mu \in \mathbb{R}, \|\varphi\|=1} \|AR(\mu)\varphi\|^2 \cdot |\Im(\mu)| < \infty. \end{aligned}$$

See also [9 Theorem XIII.25] for further characterizations of *H*-smoothness. Recall that if

$$H = H_0 + \sum_{j=1}^J B_j^* A_j$$

and *A<sub>j</sub>* are *H<sub>0</sub>*-smooth and *B<sub>j</sub>* are *H*-smooth, then the  $\Omega_{\pm}$  exist and are unitary [9 Theorem XIII.24]. This assumption on  $V = \sum_{j=1}^J B_j^* A_j$  leads to a boring situation as it rules out  $\sigma_{pp}(H)$  (if  $\sigma(H_0)$  is purely ac, e.g.), so we will relax it now.

**Definition** (Local *H*-smoothness). *A* is called *H*-smooth on a Borel set  $I \subseteq \mathbb{R}$  if  $AP_I$  is *H*-smooth. (Here  $P_I$  denotes the PVM associated to *H*.)

**Theorem XIII.30** (Sufficient criteria for local *H*-smoothness). Suppose that either

- (1)  $\sup_{\lambda \in I} \sup_{0 < |\varepsilon| < 1} |\varepsilon| \|AR(\lambda + i\varepsilon)\|^2 < \infty$  or
- (2)  $\sup_{\lambda \in I} \sup_{0 < |\varepsilon| < 1} \|AR(\lambda + i\varepsilon)A^*\| < \infty$

hold. Then *A* is *H*-smooth on  $\bar{I}$ .

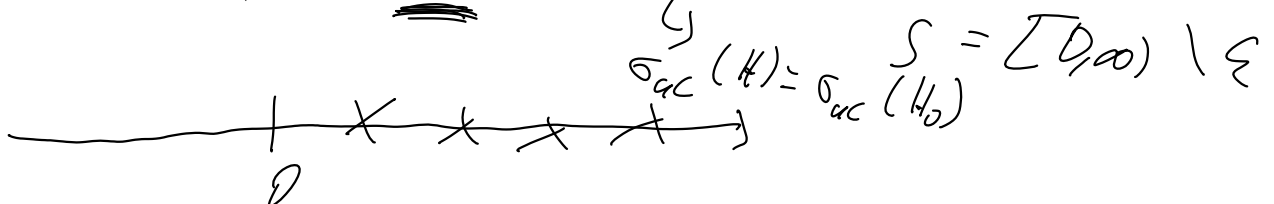
Define the local wave operators

$$W_{\pm} = s - \lim_{t \rightarrow \mp\infty} e^{itH} e^{-itH_0} P_I^{(0)}, \quad \tilde{W}_{\pm} = s - \lim_{t \rightarrow \mp\infty} e^{itH_0} e^{-itH} P_I^{(V)} = \frac{1}{2i\varepsilon} (R_V(d-i\varepsilon) - R_V(d+i\varepsilon))$$

**Theorem XIII.31** (Existence and completeness of local wave operators). If  $H = H_0 + \sum B_j^* A_j$  and *A<sub>j</sub>* are *H<sub>0</sub>*-smooth and *B<sub>j</sub>* are *H*-smooth on some open interval  $I \subseteq \mathbb{R}$ , then the local wave operators  $W_{\pm}$  and  $\tilde{W}_{\pm}$  exist and satisfy

$$W_{\pm}^* = \tilde{W}_{\pm}, \quad \tilde{W}_{\pm} W_{\pm} = P_I^{(0)}, \quad W_{\pm} \tilde{W}_{\pm} = P_I^{(V)}.$$

→ **Corollary** (Paving large sets). Let  $S \subseteq \mathbb{R}$  with  $S = \bigcup_{\ell \geq 1} I_{\ell}$  where  $I_{\ell}$  are open bounded intervals and  $H = H_0 + \sum_j B_j^* A_j$ . Suppose *A<sub>j</sub>* are *H<sub>0</sub>*-smooth and *B<sub>j</sub>* are *H*-smooth on  $I_{\ell}$  and that  $\sigma(H_0) \setminus S$  and  $\sigma(H) \setminus S$  have zero Lebesgue measure. Then the generalized wave operators  $s - \lim_{t \rightarrow \mp\infty} e^{itH} e^{-itH_0} P_{ac}^{(0)}$  exist and are complete.



## 3. DISTORTED FOURIER TRANSFORM

Let

$$H = P_0(D) + V(x, D) \quad \text{in } L^2(\mathbb{R}^d)$$

where  $P_0$  is real and simply characteristic<sup>6</sup> (see Hörmander [3], Definition 14.3.1]),  $\sigma_{pp}(P_0) = \{0\}$ , and  $V(x, D)$  is a symmetric short range perturbation of  $P_0$  in the sense of Hörmander [3], Definition 14.4.1]. Recall the Agmon–Hörmander spaces  $B$  and  $B^*$  and let

$$\begin{aligned} Z(P_0) &:= \{\lambda \in \mathbb{R} : P_0(\xi) = \lambda \text{ and } dP_0(\xi) = 0 \text{ for some } \xi \in \mathbb{R}^d\} \quad \text{and} \\ S_\lambda &:= \{\xi \in \mathbb{R}^d : P_0(\xi) = \lambda\}. \end{aligned}$$

Recall

$$\begin{aligned} \int \mathbf{1}_\Omega(\lambda)(dE_\lambda^{(0)} f, f) &= \pm \lim_{\varepsilon \searrow 0} \frac{1}{\pi} \int_{\mathbb{R}} \mathbf{1}_\Omega(\lambda) \Im(R_0(\lambda \pm i\varepsilon) f, f) d\lambda \\ &= \int_{\mathbb{R}} d\lambda \mathbf{1}_\Omega(\lambda) \int_{S_\lambda} |\hat{f}(\xi)|^2 d\sigma_{S_\lambda}(\xi), \quad f \in L^2 \end{aligned}$$

and the resolvent formula  $R(\lambda \pm i0)f = R_0(\lambda \pm i0)f_{\lambda \pm i0}$  where  $f_z = (1 + VR_0(z))^{-1}f$  is a continuous function of  $z \in \mathbb{C}^\pm \setminus (\sigma_{pp}(H) \cup Z(P_0))$  with values in  $B$ . Thus, we have

$$\int \mathbf{1}_\Omega(\lambda)(dE_\lambda^{(V)} f, f) = \int_{\mathbb{R}} d\lambda \mathbf{1}_\Omega(\lambda) \int_{S_\lambda} |\hat{f}_{\lambda \pm i0}(\xi)|^2 d\sigma_{S_\lambda}(\xi), \quad f \in B,$$

whenever  $\Omega \cap (\sigma_{pp}(H) \cup Z(P_0)) = \emptyset$ . This motivates

**Definition 1.** *If  $f \in B$ , then the  $L^2$  functions defined by*

$$\begin{aligned} (\mathcal{F}_\pm f)(\xi) &= \mathcal{F}[(1 + VR_0(\lambda \pm i0))^{-1}f](\xi), \quad \xi \in S_\lambda \\ &= \mathcal{F}[(1 - VR(\lambda \pm i0))f](\xi) \end{aligned} \tag{3.1}$$

*almost everywhere in  $S_\lambda$  are called distorted Fourier transforms of  $f$ .*

We recall the following properties of solutions of scattering states. Let  $B_{P_0}^* = \{u : P_0^{(\alpha)}u \in B^* \text{ for every } \alpha\}$ .

**Lemma 2** (Hörmander [3], Lemma 14.6.6). *If  $u \in B_{P_0}^*$ ,  $\lambda \notin Z(P_0)$ , and  $(P_0(D) + V - \lambda)u = 0$ , then  $u$  is given by the solution of the Lippmann–Schwinger equation*

$$u = u_\pm - R_0(\lambda \mp i0)Vu \tag{3.2}$$

$$= (1 - R(\lambda \mp i0)V)u_\pm, \tag{3.3}$$

where

$$\hat{u}_\pm = v_\pm \delta(P_0 - \lambda) = v_\pm d\sigma_{S_\lambda}(\xi), \quad v_\pm \in L^2(S_\lambda, d\Sigma_{S_\lambda})$$

---

<sup>6</sup>Examples are hypoelliptic operators (operators whose fundamental solutions  $E$  have  $\text{sing supp}(E) = \{0\}$ ) or operators of real principal type (operators whose principal symbol  $P_m$  is real and  $P'_m(\xi) \neq 0$  for  $\xi \neq 0$ )

and

$$\int_{S_\lambda} (|v_+|^2 - |v_-|^2) d\sigma_{S_\lambda}(\xi) = 0 \quad (3.4)$$

where  $d\sigma_{S_\lambda}(\xi) = |\nabla P_0(\xi)|^{-1} d\Sigma_{S_\lambda}(\xi)$  and  $d\Sigma_{S_\lambda}(\xi)$  is the euclidean surface measure on  $S_\lambda$ . Moreover, if  $\lambda \notin (Z(P_0) \cup \sigma_{pp}(P_0 + V))$ , then

$$(\mathcal{F}_+ f, \hat{u}_+) = (\mathcal{F}_-, \hat{u}_-) = (f, u), \quad \text{if } f \in B. \quad (3.5)$$

Let us also recall

**Theorem 3** (Hörmander [3] Lemma 14.6.4 and Theorem 14.6.5].  $\mathcal{F}_\pm : E^c L^2(\mathbb{R}^d) \rightarrow \widehat{L^2(\mathbb{R}^d)}$  is an isometric operator, which vanishes on  $E^{pp} L^2(\mathbb{R}^d)$ , with

$$\|E^c f\|_2^2 = \int_{\mathbb{R}^d} |\mathcal{F}_\pm f(\xi)|^2 d\xi.$$

Moreover, the intertwining property

$$\mathcal{F}_\pm e^{itH} = e^{itP_0(\xi)} \mathcal{F}_\pm$$

holds for all  $t \in \mathbb{R}$ . In particular, the restriction of  $H$  to  $E^c L^2$  is absolutely continuous (since  $P_0$  has purely absolutely continuous spectrum).

Moreover,  $\mathcal{F}_\pm : E^c L^2(\mathbb{R}^d) \rightarrow \widehat{L^2(\mathbb{R}^d)}$  is actually unitary, i.e., the restriction of  $H$  to  $E^c L^2$  is unitarily equivalent to  $P_0$ , i.e.,  $\sigma_c(H) = \sigma_{ac}(H) = \sigma(P_0)$ . In particular, for  $f \in E^c(L^2(\mathbb{R}^d))$ , we have

$$(\mathcal{F}_\pm H f)(\xi) = P_0(\xi)(\mathcal{F}_\pm f)(\xi), \quad \text{i.e.,} \quad (H f)(x) = (\mathcal{F}_\pm^* P_0(\cdot) \mathcal{F}_\pm f)(x).$$

In particular, it follows that

$$\mathcal{F}_\pm^* \mathcal{F}_\pm = E^c \quad \text{and} \quad \mathcal{F}_\pm \mathcal{F}_\pm^* = \mathbf{1}_{\widehat{L^2}}.$$

The distorted Fourier transform (3.1) can be conveniently represented using the solutions  $\varphi_{\xi(\lambda)}(x)$  (for  $\xi(\lambda) \in S_\lambda$ ) of the Lippmann–Schwinger equation (3.2). In fact, we have (see also Ikebe [5] and Yafaev [11] Sections 6.6-6.8)

$$(\mathcal{F}_\pm f)(\xi) = \langle \varphi_\xi, f \rangle, \quad \xi \in \bigcup_{\lambda \in \sigma_{ac}(H)} S_\lambda \quad (3.6)$$

$$(\mathcal{F}_\pm^* g)(x) = \int_{\mathbb{R}^d} \varphi_\xi(x) g(\xi) d\xi = \int_{\sigma_{ac}(H)} d\lambda \int_{S_\lambda} d\sigma_{S_\lambda}(\xi) \varphi_\xi(x) g(\xi). \quad (3.7)$$

Moreover, we have the following expansion theorem (see also Ikebe [5] Theorem 5])

$$f = \sum_{\lambda \in \sigma_{pp}(H)} |\psi_\lambda\rangle \langle \psi_\lambda, f \rangle + \int_{\mathbb{R}^d} |\varphi_\xi\rangle \langle \varphi_\xi, f \rangle d\xi \quad (3.8)$$

where  $\{\psi_\lambda\}_{\lambda \in \sigma_{pp}(H)}$  denote the  $L^2$ -normalized eigenfunctions of  $H$ , i.e.,  $H\psi_\lambda = \lambda\psi_\lambda$ . Moreover,

$$H f = \sum_{\lambda \in \sigma_{pp}(H)} \lambda |\psi_\lambda\rangle \langle \psi_\lambda, f \rangle + \int_{\mathbb{R}^d} P_0(\xi) |\varphi_\xi\rangle \langle \varphi_\xi, f \rangle d\xi. \quad (3.9)$$

The above results motivate in particular the following definition of the *distorted Fourier restriction and extension operators*

$$(F_{S_\lambda} f)(\xi) = \langle \varphi_\xi, f \rangle = (\mathcal{F}_\pm f)(\xi), \quad \xi \in S_\lambda \quad (3.10)$$

$$(F_{S_\lambda}^* g)(x) = \int_{S_\lambda} d\sigma_{S_\lambda}(\xi) \varphi_{\xi(\lambda)}(x) g(\xi) \quad (3.11)$$

which are defined with respect to the canonical measure  $d\sigma_{S_\lambda}$ . In particular, we have for any  $\Lambda \subseteq \sigma_{ac}(H)$ ,

$$E_H(\Lambda) = \int_{P_0^{-1}(\Lambda)} |\varphi_\xi\rangle \langle \varphi_\xi| d\xi = \int_\Lambda d\lambda \int_{S_\lambda} d\sigma_{S_\lambda}(\xi) |\varphi_{\xi(\lambda)}\rangle \langle \varphi_{\xi(\lambda)}| = \int_\Lambda d\lambda F_{S_\lambda}^* F_{S_\lambda}$$

in a suitable weak sense and in particular, for  $\lambda \in \sigma_{ac}(H)$ ,

$$\frac{dE_H(\lambda)}{d\lambda} = \int_{S_\lambda} d\sigma_{S_\lambda}(\xi) |\varphi_{\xi(\lambda)}\rangle \langle \varphi_{\xi(\lambda)}| = F_{S_\lambda}^* F_{S_\lambda}.$$

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